

# Computer Animation

## 11-Kinematics and Physics

### SS 19

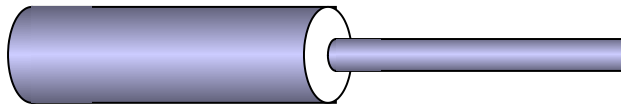
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# Hierarchical modeling

- Hierarchical modeling is placing constraints on objects organized in a tree like structure
- Examples can be:
  - A planet system
  - A robot arm
- The latter is quite common in graphics: it is constituted by objects connected end to end to form a multibody jointed chain
- These are called *articulated figures*
- They stem from robotics
- Robotics literature speaks with a different terminology:
  - Manipulator: the sequence of objects connected by joints
  - Links: the rigid objects making the chain
  - Effector: the free end of the chain
  - Frame: local coordinate system associated to each link

# Hierarchical modeling

- In graphics, most of the links are revolute joints: here one link rotates around a fixed point of the other link
- The other interesting joint for graphics is the prismatic joint, where one link translates relative to the other
- Joints restrain the degree of freedom (DOF) of the links
- Joints with more than one degree of freedom are called *complex*
- Typically, when a joint has  $n > 1$  DOF it is modeled as a set of  $n$  one degree of freedom joints



# Hierarchical modeling

- Humans and animals can be modeled as hierarchical linkages
- These are represented as a tree structure of nodes connected by arcs
- The highest node of this structure is called the root node, and is the node that has position WRT the global coordinate system
- All other nodes have their position only as relative to the root node
- A node that has no child is called a leaf node
- Each node contains the info necessary to define the position of the corresponding part
- Two types of transformations are associated with an arc leading to a node:
  - Rotation and translation of the object to its position of attachment to the father link
  - Information responsible for the joint articulation

# Hierarchical modeling

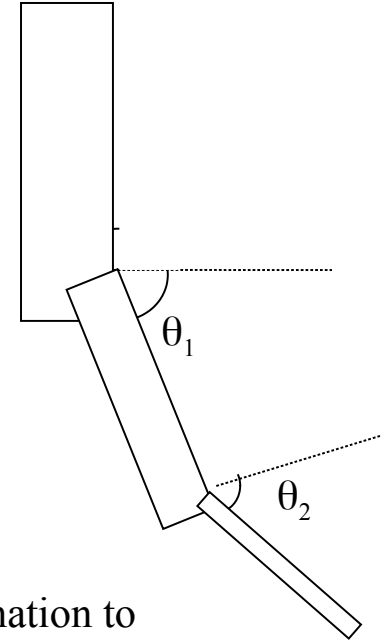
- How does this work?
- The idea is simple, store at each node
  - Info on the node geometry
  - The transformation (its rotation) with respect to the father node in the tree
- To obtain the position of the  $i$ -th node in the chain, one has to simply multiply the transformations to obtain the position of the current arc to be displayed
- The root node of course contains info of its absolute position and orientation in the global coord. system

$T_0$ : transformation to rotate  $K_0$  in WCS

$T_1$ : transformation to rotate  $K_1$  WRT  $K_0$   
= rotation by  $\theta_1$

$T_2$ : transformation to rotate  $K_2$  WRT  $K_1$   
= rotation by  $\theta_2$

- To obtain the position of  $K_2$  in WCS, one will then have to multiply  $T_0 T_1 T_2$



# Forward kinematics

- Traversing the tree of the nodes produces the correct picture of the object
- Traversal is done depth first until a leaf is met
- Once the corresponding arc is evaluated, the tree is backtracked up until the first unexplored node is met
- This is repeated until there are no nodes left unexplored
- A stack of transforms is kept
- When tree is traversed down-wards, the corresponding transformation is added to the stack
- Moving up pops the transformation from the stack
- Current node position is generated through multiplying the current stack transforms

# Forward kinematics

- To animate the whole, the rotation parameters are manipulated and the corresponding transforms are actualized
- A complete set of rotations on the whole arcs is called a *pose*
- A pose is obviously a vector of rotations
- Moving an object by positioning all its single arcs manually is called forward kinematics
- This is not so user-friendly
- Instead of specifying the whole links, the animator might want to specify the end position of the effector
- The computer computes then the position of the other links
- This is called *inverse kinematics*

# Inverse kinematics

- The user gives the position of the end effector and the computer computes the joint angles
- One can have zero, one or multiple solutions
  - No solution: overconstrained problem
  - Multiple solutions: underconstrained problem
  - Reachable workspace: volume that end effector can reach
  - Dextrous workspace: volume that end effector can reach in any orientation
- Computing the solution to the problem can at times be tricky
- If the mechanism is simple enough, then the solution can be computed analytically
- Given an initial and a final pose vector, the solution can be computed by interpolating the values of the pose vector
- If the solution cannot be computed analytically, then there is a method based on the jacobian to compute incrementally a solution

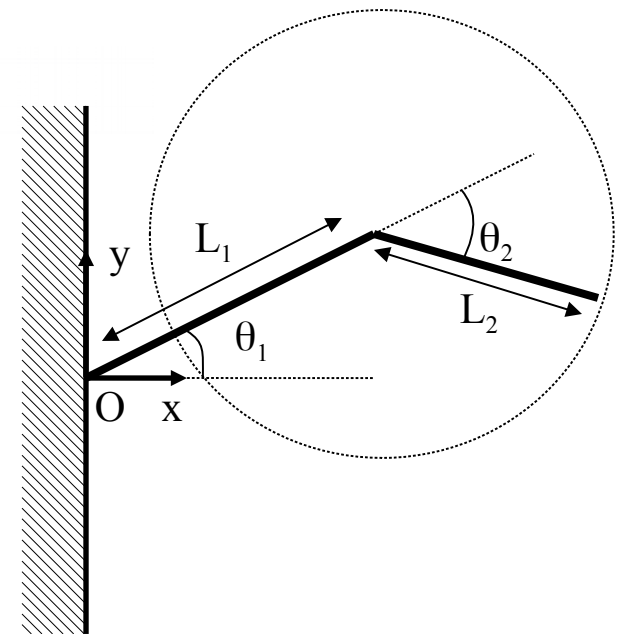


# Inverse kinematics

- Consider the figure: the 2<sup>nd</sup> arm rotates around the end of the 1<sup>st</sup> arm.
- It is clear that all positions between  $|L_1 - L_2|$  and  $|L_1 + L_2|$  can be reached by the arm.
- Set the origin like in the drawing
- In inverse kinematics, the user gives the  $(X, Y)$  position of the end effector

- Obviously there are only solutions if

$$|L_1 - L_2| \leq \sqrt{X^2 + Y^2} \leq |L_1 + L_2|$$



# Inverse kinematics

- $\cos\theta_T = X/(X^2+Y^2)^{1/2}$   
 $\Rightarrow \theta_T = \arccos(X/(X^2+Y^2)^{1/2})$
- Because of the cosine rule we have also that

$$\cos(\theta_1 - \theta_T) =$$

$$(L_1^2 + X^2 + Y^2 - L_2^2) / 2L_1(X^2 + Y^2)^{1/2}$$

and

$$\cos(\pi - \theta_2) =$$

$$(L_1^2 + L_2^2 - (X^2 + Y^2)) / 2L_1L_2$$

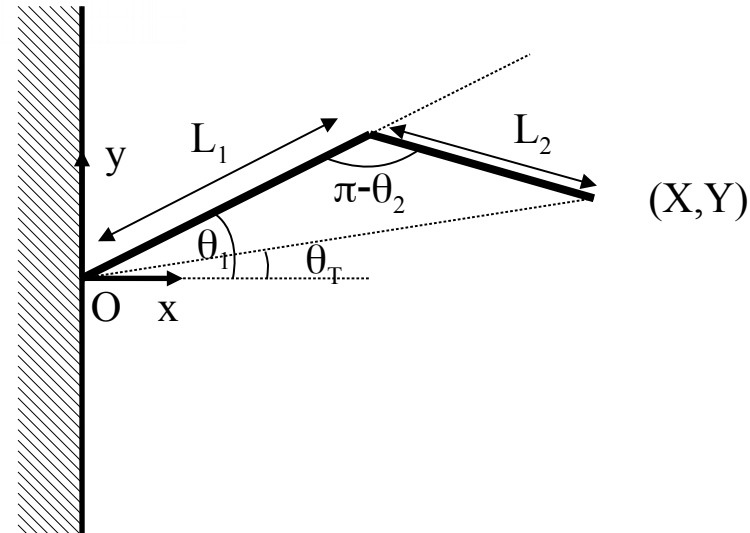
from which we have

$$\theta_1 = \arccos\left(\frac{(L_1^2 + X^2 + Y^2 - L_2^2)}{2L_1(X^2 + Y^2)^{1/2}}\right) + \theta_T$$

and

$$\theta_2 = \arccos\left(\frac{(L_1^2 + L_2^2 - (X^2 + Y^2))}{2L_1L_2}\right)$$

- Note that two solutions are possible, symmetric with respect to the line joining the origin and  $(X, Y)$



# Inverse kinematics

- In general, for the quite simple armatures used in robotics it is possible to implement such analytic solutions
- Unfortunately this works only for simple cases
- For more complicated armatures, the number of possible solutions there may be infinite solutions for a given effector location, and computations become so difficult to do that iterative numeric solution must be used

# Using the Jacobian

- When the solution is not analytically computable, incremental methods converging to the solution are used
- To do this, the matrix of the partial derivatives has to be computed
- This is called the *Jacobian*

- Suppose you have six independent variables and you have a six unknowns that are functions of these variables

$$y_1 = f_1(x_1, x_2, x_3, x_4, x_5, x_6)$$

$$y_2 = f_2(x_1, x_2, x_3, x_4, x_5, x_6)$$

$$y_3 = f_3(x_1, x_2, x_3, x_4, x_5, x_6)$$

$$y_4 = f_4(x_1, x_2, x_3, x_4, x_5, x_6)$$

$$y_5 = f_5(x_1, x_2, x_3, x_4, x_5, x_6)$$

$$y_6 = f_6(x_1, x_2, x_3, x_4, x_5, x_6)$$

or, in vector notation,

$$\underline{Y} = \underline{F}(\underline{X})$$

# Using the Jacobian

- What happens when the input variables change?
- The equations can be written in differential form:

$$\begin{aligned} \delta y_i = & \frac{\partial f_i}{\partial x_1} \delta x_1 + \frac{\partial f_i}{\partial x_2} \delta x_2 \\ & + \frac{\partial f_i}{\partial x_3} \delta x_3 + \frac{\partial f_i}{\partial x_4} \delta x_4 \\ & + \frac{\partial f_i}{\partial x_5} \delta x_5 + \frac{\partial f_i}{\partial x_6} \delta x_6 \end{aligned}$$

or, in vector form

$$\delta \underline{Y} = \frac{\partial \underline{F}}{\partial \underline{X}} \delta \underline{X}$$

- Given n equations in n variables, the matrix

$$J = \begin{bmatrix} \frac{\partial f_1}{\partial x_1} & \frac{\partial f_1}{\partial x_2} & \dots & \frac{\partial f_1}{\partial x_n} \\ \frac{\partial f_2}{\partial x_1} & \frac{\partial f_2}{\partial x_2} & \dots & \frac{\partial f_2}{\partial x_n} \\ \vdots & \vdots & \ddots & \vdots \\ \frac{\partial f_n}{\partial x_1} & \frac{\partial f_n}{\partial x_2} & \dots & \frac{\partial f_n}{\partial x_n} \end{bmatrix}$$

is called the Jacobian matrix of the system

- The Jacobian can be seen as a mapping of the velocities of  $\underline{X}$  to velocities of  $\underline{Y}$

# Summary: articulated bodies

- Very useful for enforcing certain relationships among elements of an animation
- Allows animator to concentrate on effector forgetting the rest of the body
- Damn hard to do, to date not real in real time
- Adding control expressions can be tricky
- No physics considered. Only kinematics

# Physics in animation

- While animating through kinematics may be interesting for plenty of applications, integrating physics is more difficult and a challenging problem
- The easiest way of integrating physics is rigid body simulation
- While physics is concerned with the exactness of the representation, animation is more interested in „credible“ effects, and in rendering frame by frame
- Having to deal with the system at discrete time samples creates numerical problems in the solution methods which are not simple to deal with

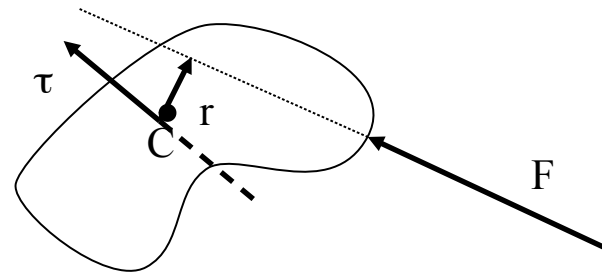
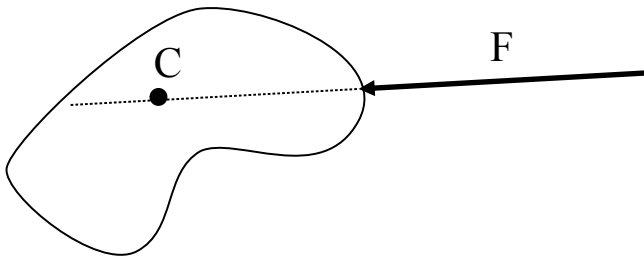
# Recap on physics (physics 101)

- In the equations of motion, the following quantities play a role
  - Distance = speed · time
  - time = frame# · timeperframe
  - averageVelocity = distance traveled / time
- Linear motion U
  - s = position
  - v = velocity
  - a = acceleration
  - $s(t) = v(t) \cdot t$
  - $v(t) = a(t) \cdot t$
  - $s(t) = \frac{1}{2} a(t) \cdot t^2$
- Circular motion
  - $\theta$  angular position
  - $\omega$  angular velocity
  - $\theta(t) = \omega(t) \cdot t$
  - For a body in circular motion, we have  $a(t) = (-\omega)^2 \cdot r$
- Newton's law:
  - $\underline{F} = m \cdot \underline{a}$
  - A body continues its own motion therefore if the sum of the forces acting on it = 0
  - $\Sigma F_i = 0$




# Recap on physics (physics 101)

- Remember the definition of center of mass
  - The point at which the object is balanced in all directions
  - If an external force is applied to a body in line with its center of mass, then the body would move as if it was a point at the center of mass C
- Torque is the tendency of a force to produce circular motion
  - It is produced by a force off center to the center of mass
  - $\tau = r \times F$
  - Clearly,  $\tau \perp F$  and  $\tau \perp r$
- An object does not move if  $\Sigma F_i = 0$  and  $\Sigma \tau_i = 0$



# Recap on physics (physics 101)

- Linear springs:
  - Hooke's law:  
$$F = -k \cdot x,$$
where  $x$  is the change from the equilibrium length of the spring
- Friction:
  - Static:  
$$F_s = s \cdot f_N$$
where  $F_s$  = frictional force  
 $s$  = static friction coefficient   
 $f_N$  = normal force
  - Kinetic:  
$$F_k = k \cdot f_N$$
with similar coefficient definitions as in static friction
- Momentum:  $m \cdot v$ 
  - In a closed system, total momentum does not vary
- Angular momentum:  
$$L = r \times p$$
where  
 $r$  = vector from center of rotation  
 $p$  = momentum ( $m \cdot v$ )
  - Note that  
$$\tau = dL/dt$$
  - In a closed system, total angular momentum does not vary
- Inertia tensor: the resistance of an object to change its angular momentum

# Rigid body simulation

- If one wants to simulate rigid bodies, many forces act on them
- Such forces vary in time continuously and in a non linear way
- Therefore it is not enough to evaluate velocities and accelerations at fixed timesteps  $\Delta t$
- Evaluating the velocities at  $t_0$ ,  $t_0+\Delta t$ ,  $t_0+2\Delta t$  does not generate a correct movement, and slowly drifts away from the correct solution
- This solution method is an example of the Euler integration method
- The accuracy of the method is determined by the size of the time step
- Obviously the shorter the time step, the more computations are needed
- A better way of integrating the equations bases on the Runge Kutta method
  - In particular, often 2nd order Runge Kutta (midpoint method) is used
  - Remember, the order of the RK method is the magnitude of the error term
  - Even higher order ones, 4th or 5th ones are used

# Motion equations for a rigid body

- To develop the equation of motion for a rigid body, we have to apply some of the physics presented before
- When a force is applied to a rigid body, the force and the relative torque are applied to the body
- To uniquely solve for the resulting motions of interacting bodies, linear and angular momentum have to be conserved
- Finally, to calculate the angular momentum the distribution of an object mass in space has to be characterized with its inertia tensor.

# Orientation and rotational movement

- Similar to position, velocity and acceleration, 3D objects have
  - orientation,
  - angular velocity and
  - angular acceleration
- which vary in time
- Let  $R(t)$  represent the object rotation
- *Angular velocity*  $\underline{\omega}(t)$  is the rate at which the object is rotated (independent from linear velocity)
- The direction of  $\underline{\omega}(t)$  indicates the orientation of the axis about which the object is rotating
- The magnitude of  $\underline{\omega}(t)$  gives the speed of rotation in revs per unit time

# Center of mass

- The center of mass of a body is defined as the integral of the differential mass times its position in the object
- In a body with discrete masses, then the center of mass is *at*  $q_i(t)$ , the center of mass is at  $x(t) = \sum m_i q_i(t) / \sum m_i$

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# Forces and torque

- A linear force applied to a mass gives rise to a linear acceleration

$$F=ma \quad (\text{Newton's law})$$

- The various forces applied to a point sum up

$$F(t)=\sum f_i(t)$$

- The torque arising from the application of forces acting on a point of an object is given by

$$\tau_i(t)=(q(t)-x(t)) \times f_i(t)$$

$$\tau(t)=\sum \tau_i(t)$$

# Momentum

- The momentum of an object (= mass x velocity) is decomposed into
  - linear component: acts on center of mass
  - angular components: acts WRT center
- Both are preserved in a closed system
- Linear momentum  $p = m v$
- Total linear momentum of a rigid body:  $P(t) = \sum m_i \dot{q}_i(t)$
- Deriving  $p = mv$  we obtain  $P^\circ(t) = M v^\circ(t) = F(t)$
- Angular momentum is a measure of the rotating mass weighted by the mass's distance from the axis of rotation
- $L(t) = \sum ((q(t) - x(t)) \times m_i (\dot{q}^\circ(t) - v(t)))$   
 $= \sum (R(t)q \times m_i (\omega(t) \times (q(t) - x(t))))$   
 $= \sum (m_i (r(t)q \times (\omega(\tau) \times R(t)q)))$
- Similar to linear momentum, torque equals the change in angular momentum  
 $L^\circ(t) = \tau(t)$
- Note that since angular momentum depends on distance to center of mass, to maintain constant angular momentum, the angular velocity increases if the distance of the mass decreases



# Inertia tensor

- Angular momentum is related to angular velocity the same way linear momentum is related to linear velocity  $P(t) = M \cdot v(t)$
- We have  $L(t) = I(t) \cdot \omega(t)$
- The distrib. of mass of the obj. in space is defined through a matrix, the inertia tensor  $I(t)$

$$I_{obj} = \begin{bmatrix} I_{xx} & I_{xy} & I_{xz} \\ I_{yx} & I_{yy} & I_{yz} \\ I_{zx} & I_{zy} & I_{zz} \end{bmatrix}$$

where the matrix terms are computed by integrating over the object, and  $I$  is symmetric

- In general,  $I_{xx} = \iiint \rho(q)(q_y^2 + q_z^2) dx dy dz$  where  $\rho$  is the density of at an obj. point  $q = (q_x, q_y, q_z)$
- In the case of discrete masses  $I_{xx} = \sum m_i (y_i^2 + z_i^2)$ ,  $I_{xy} = \sum m_i x_i y_i$   
 $I_{yy} = \sum m_i (x_i^2 + z_i^2)$ ,  $I_{xz} = \sum m_i x_i z_i$   
 $I_{zz} = \sum m_i (x_i^2 + y_i^2)$ ,  $I_{yz} = \sum m_i y_i z_i$
- In a center of mass centered obj space, the inertia tensor of a transformed object depends on the obj orientation but not on its position, and therefore it depends on time
- It can be transformed with  $I(t) = R(t) I_{obj} R(t)^T$

# Motion equations

- The state of an object can be determined by the vector containing
  - Position
  - Orientation
  - Linear momentum
  - Angular momentum

$$S(t) = \begin{bmatrix} x(t) \\ R(t) \\ P(t) \\ L(t) \end{bmatrix}$$

- Object mass and its object space inertia tensor  $I_{\text{obj}}$  do not change in time

- At any time, the following quantities can be computed:
  - Inertia tensor  $I(t) = R(t) I_{\text{obj}} R(t)^T$
  - Angular vel.  $\omega(t) = I(t)^{-1} L(t)$
  - Linear vel.  $v(t) = P(t) / M$

- Now the time derivative can be formed:

$$\frac{d}{dt} S(t) = \frac{d}{dt} \begin{bmatrix} x(t) \\ R(t) \\ P(t) \\ L(t) \end{bmatrix} = \begin{bmatrix} v(t) \\ \omega(t)^* R(t) \\ F(t) \\ \tau(t) \end{bmatrix}$$

This is enough to run a simulation

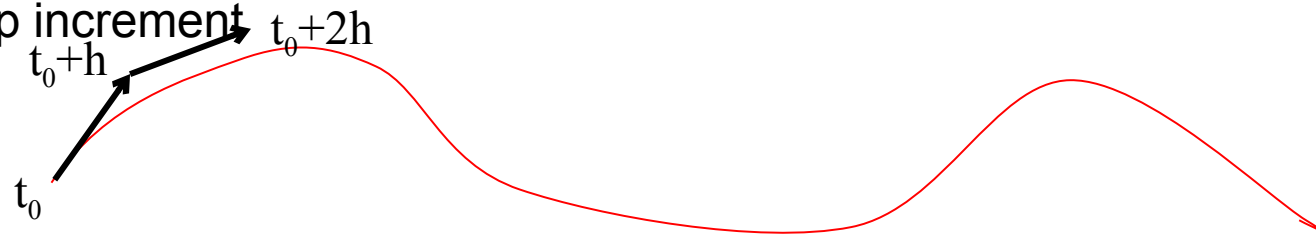
- A differential equation solver can now be used

# Motion equations

- As the simplest solver, one can use Euler's method
  - The values of the state array are updated by multiplying their time derivatives by the length of the time step
- Euler methods assume that the solution at the time point  $t_0+h$  will can be approximated by the solution at  $t_0$  plus the stepsize  $h$  multiplied by the derivative at the time  $t_0$ :

$$y(t_0+h) \approx y(t_0) + hf'(x(t_0), y(t_0))$$

- This is equivalent to saying: for a little while ( $h$ ), approximating my movement with the tangent to the solution at the previous step is okay.
- Or in other words, since the derivative of space is velocity, it is the same as approximating the movement for a little while with the velocity times the timestep increment



- As long as the timestep  $h$  is small the resulting solutions are correct.
- There are, of course, more refined methods for solving differential equations, but explaining them would be beyond the purpose of this course.



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